



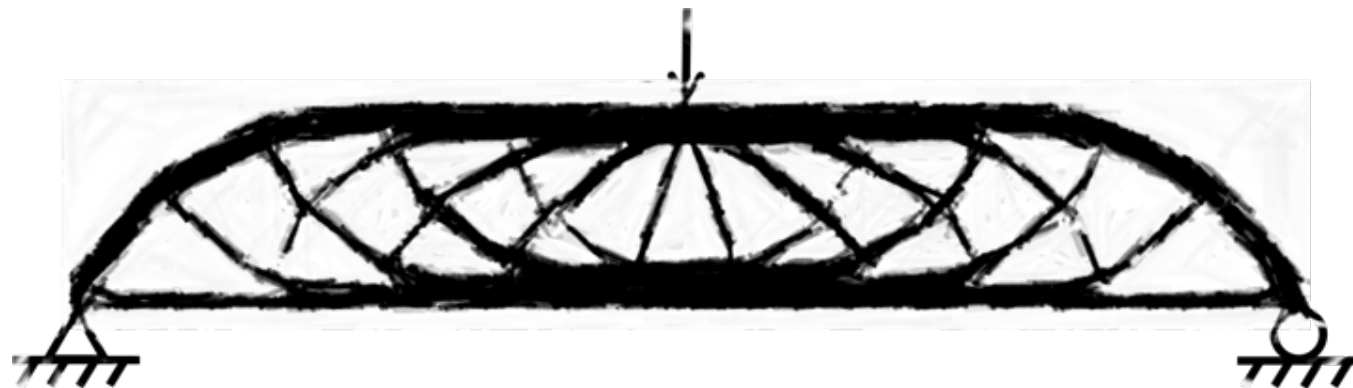
11th US National Congress on Computational Mechanics:

ON RESTRICTION METHODS FOR TWO-PHASE OPTIMAL SHAPE PROBLEMS

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- For example, we show that a consequence of the ill-posedness is that **smearing** of Heaviside function transforms the topology problem into the **variable thickness** problem



- The **two-phase** optimal shape problem is defined as:

$$\inf_{\chi \in \mathcal{A}} J(\chi, \mathbf{u}_\chi) \quad \text{where } \mathbf{u}_\chi \in \mathcal{V} \text{ solves } \mathcal{B}(\mathbf{u}, \mathbf{v}; \chi) = \ell(\mathbf{v}), \quad \forall \mathbf{v} \in \mathcal{V}$$



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$\mathcal{A} \subseteq L^\infty(\Omega; \{0, 1\})$ is the given space of **admissible** designs,

$$\mathcal{B}(\mathbf{u}, \mathbf{v}; \chi) = \int_{\Omega} \boldsymbol{\epsilon}(\mathbf{u}) : [\chi \mathbf{C}^+ + (1 - \chi) \mathbf{C}^-] : \boldsymbol{\epsilon}(\mathbf{v}) \, d\mathbf{x}, \quad \ell(\mathbf{v}) = \int_{\Gamma_N} \mathbf{t} \cdot \mathbf{v} \, ds$$

are the bilinear and linear forms, and $\mathcal{V} = \{ \mathbf{u} \in H^1(\Omega; \mathbb{R}^d) : \mathbf{u}|_{\Gamma_D} = \mathbf{0} \}$



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- The objective function $J(\chi, \mathbf{u})$ is assumed to be **continuous** in strong topology of $L^1(\Omega) \times H^1(\Omega; \mathbb{R}^d)$
 - The objective function for the **minimum compliance** is given by

$$J(\chi, \mathbf{u}_\chi) = \ell(\mathbf{u}_\chi) + \lambda \int_{\Omega} \chi \, d\mathbf{x}$$

where λ is the volume penalty parameter



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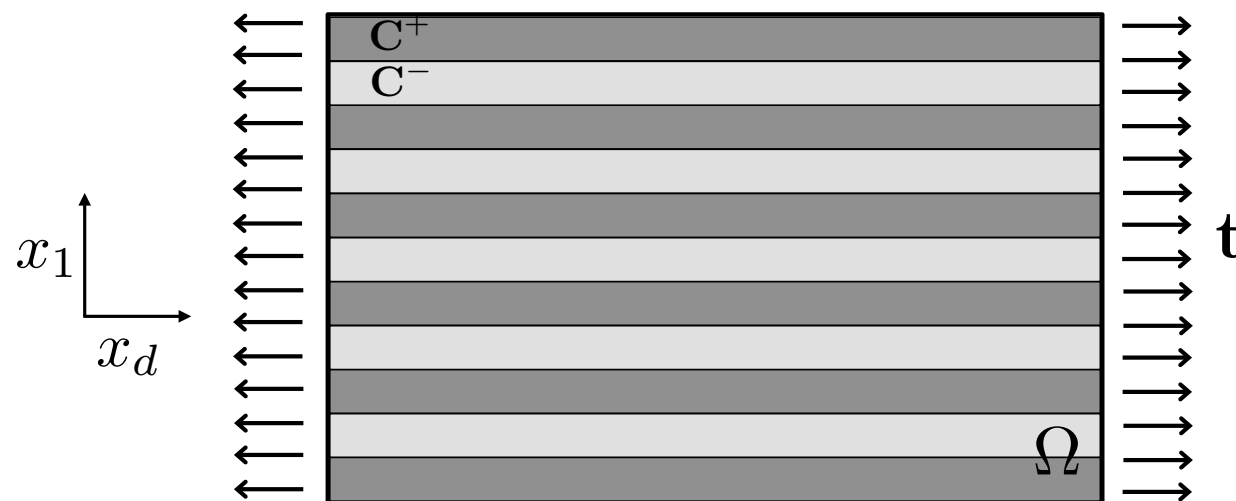
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- The optimal design for this problem is a rank-1 laminate, whose stiffness is precisely the H -limit of $\chi_n C^+ + (1 - \chi_n) C^-$



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PROPOSITION: *Let $\chi_n, \hat{\chi} \in L^\infty(\Omega; [0, 1])$ be such that $\chi_n \rightarrow \hat{\chi}$ in $L^1(\Omega)$. Then, up to a subsequence, the associated state solutions also converge, i.e., $\mathbf{u}_{\chi_n} \rightarrow \mathbf{u}_{\hat{\chi}}$ in $H^1(\Omega)$.*



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- A well-known example is the space of designs with **bounded perimeter**:

$$\mathcal{A} = \{ \chi \in BV(\Omega; \{0, 1\}) : \int_{\Omega} |\nabla \chi| \, d\mathbf{x} \leq \bar{P} \}$$



- Another choice (Liu et al. 2003) is to set $\mathcal{A} = H(\mathcal{F})$ where the implicit functions $\varphi \in \mathcal{F} \subseteq W^{1+\theta,2}$ satisfy:

$$(R1) : \quad \|\varphi\|_{W^{1+\theta,2}(\Omega)} \leq M$$

$$(R2) : \quad |\varphi(\mathbf{x})| + |\nabla\varphi(\mathbf{x})| \geq \nu \quad \text{a.e. } \mathbf{x} \in \Omega$$

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- (R1) excludes the possibility of rapid oscillations of the implicit functions:

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- (R2) ensures that the phase boundary

$$\{\mathbf{x} \in \Omega : \varphi(\mathbf{x}) = 0\},$$

which is where the Heaviside is discontinuous, has zero measure:

- Without it, $\varphi_n(\mathbf{x}) = (\alpha/n^{2+\theta}) \sin(nx_1)$ gives a minimizing sequence that satisfies (R1) but does not converge



- If no restrictions are placed on φ , the usual **approximation** of the Heaviside by

$$H_w(\varphi)(\mathbf{x}) = \begin{cases} 0, & \varphi(\mathbf{x}) < -w \\ h_w(\varphi(\mathbf{x})), & |\varphi(\mathbf{x})| \leq w \\ 1, & \varphi(\mathbf{x}) > w \end{cases}$$

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- For any $\rho \in L^\infty(\Omega; [0, 1])$, there exists $\varphi \in L^\infty(\Omega; [-\alpha, \alpha])$ such that $\rho = H_w(\varphi)$. Conversely, $H_w(\varphi)$ represents a thickness function
- Note also that the conditions of optimality are the same:

$$H'_w(\varphi) [\lambda - E(\mathbf{u})] = 0 \quad \text{when } -w < \varphi < w$$

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Therefore the optimal solution with such approximation will contain large "grey" regions filled by the intermediate phases



- (R1) can be imposed via convolution with a **smooth filter**, i.e., by defining
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$$J_\beta(\chi, \mathbf{u}_\chi) = J(\chi, \mathbf{u}_\chi) + \beta \int_{\Omega} \chi(1 - \chi) \, d\mathbf{x}$$

OR change the state equation to penalize the intermediate stiffnesses:

$$\mathcal{B}_p(\mathbf{u}, \mathbf{v}; \chi) = \int_{\Omega} \boldsymbol{\epsilon}(\mathbf{u}) : [\chi^p \mathbf{C}^+ + (1 - \chi^p) \mathbf{C}^-] : \boldsymbol{\epsilon}(\mathbf{v}) \, d\mathbf{x}$$



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In both cases, separation of phases and thus transversality is achieved in the optimal regime.

- The condition of optimality for $\{-w < \varphi < w\}$ respectively are:

$$H'_w(\varphi) \{ \lambda + \beta [1 - 2H_w(\varphi)] - E(\mathbf{u}) \} = 0$$

$$H'_w(\varphi) \{ \lambda - p [H_w(\varphi)]^{p-1} E(\mathbf{u}) \} = 0$$



- The **continuum** parameters (i.e., those independent of the mesh size) are:
 - α : bound for implicit function field
 - R : radius of filtering kernel K
 - w : width of the approximate Heaviside
 - p : parameter for penalization of intermediate stiffnesses

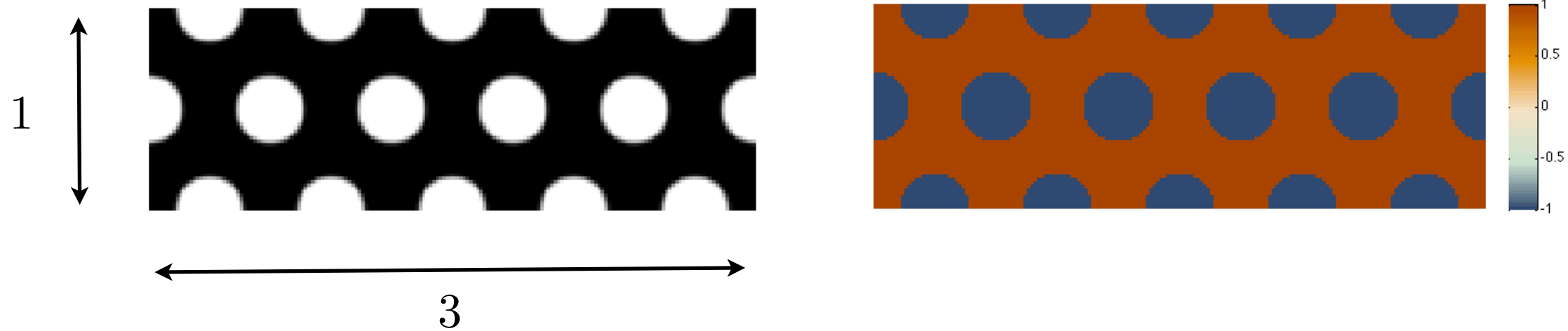
- It is not be easy to establish an **explicit** relationship between ν with above parameters in general

- However the compliance problem, the transversality constant ν is directly related to α/R (which is why we set w to be fixed fraction of α/R)

Some numerical results



Initial guess:

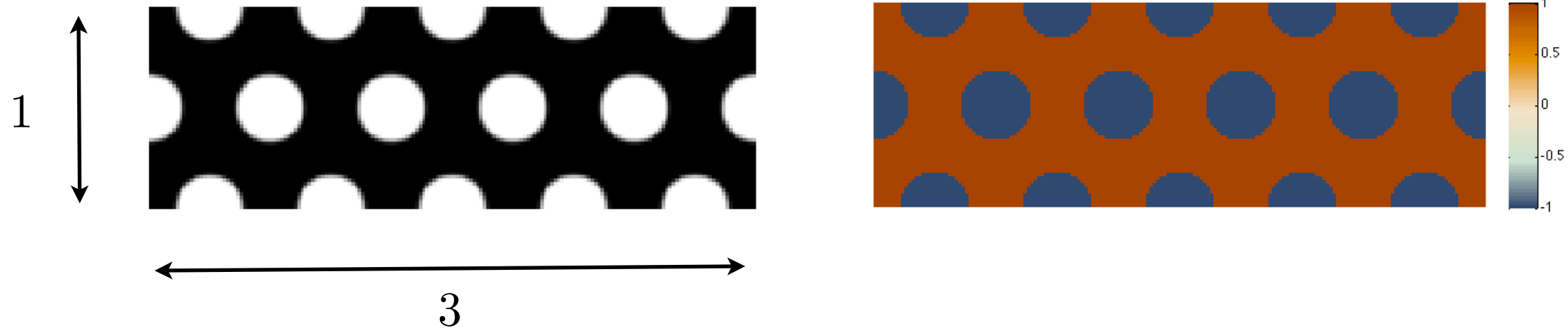


Values of parameters used: $\alpha = 1$, $w = 0.0375\alpha/R$, $p = 4$

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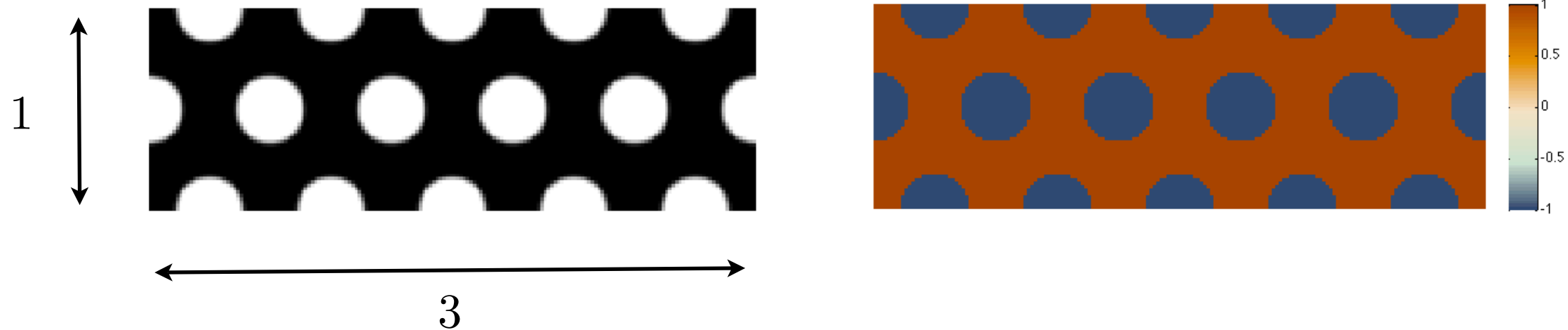
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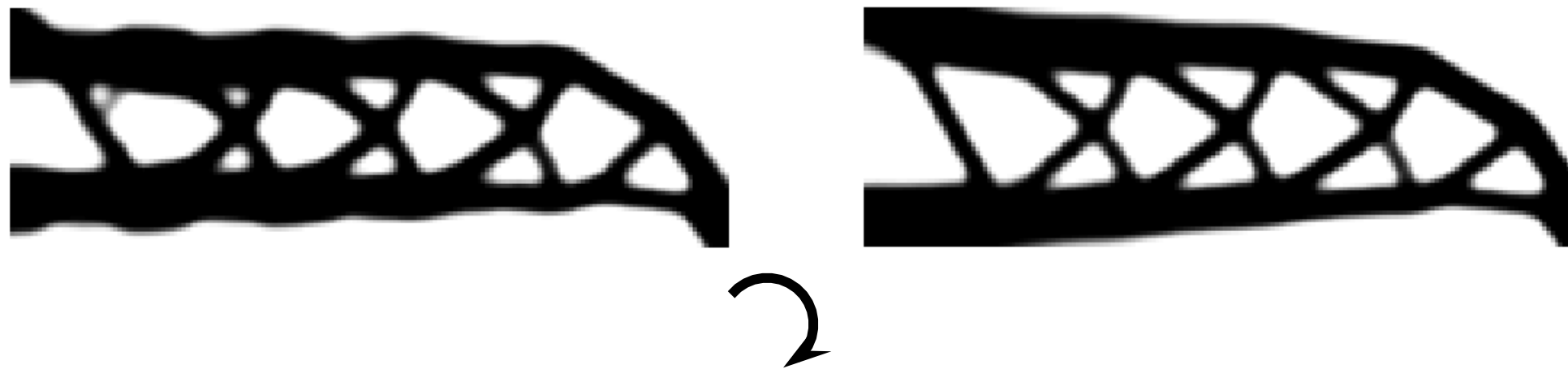
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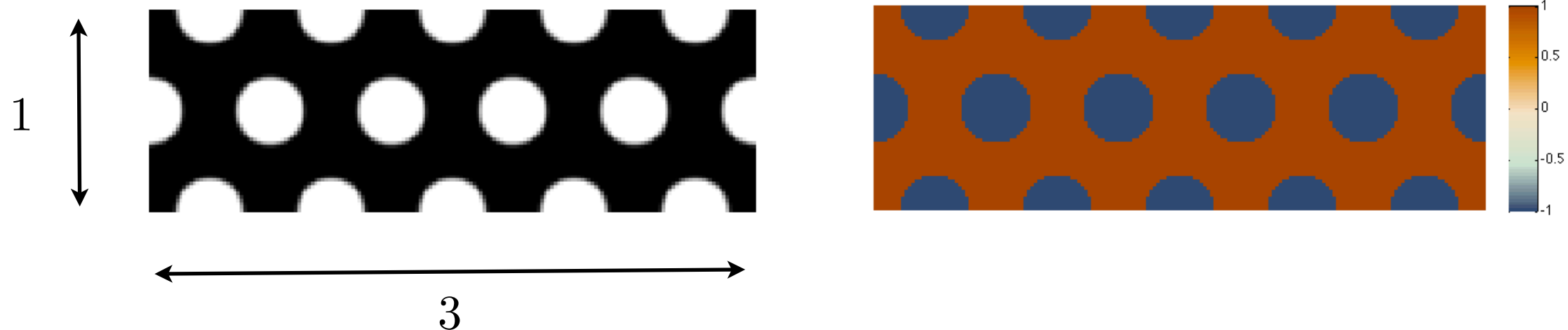
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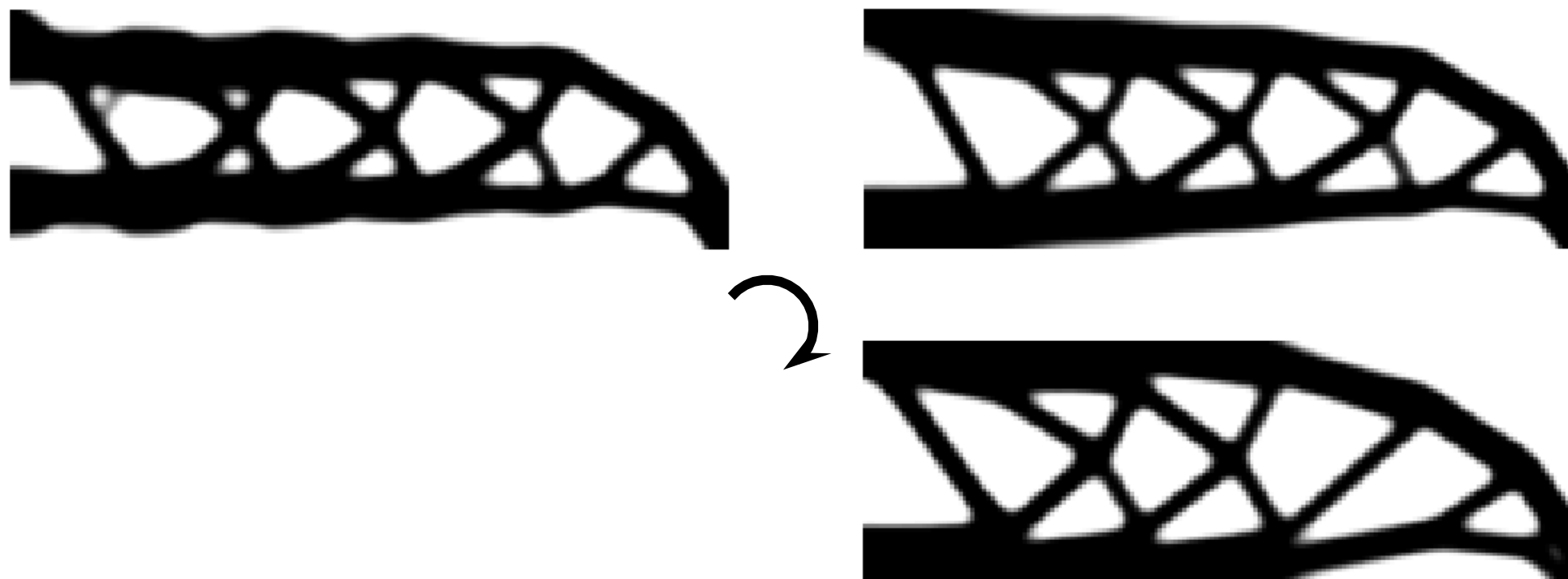
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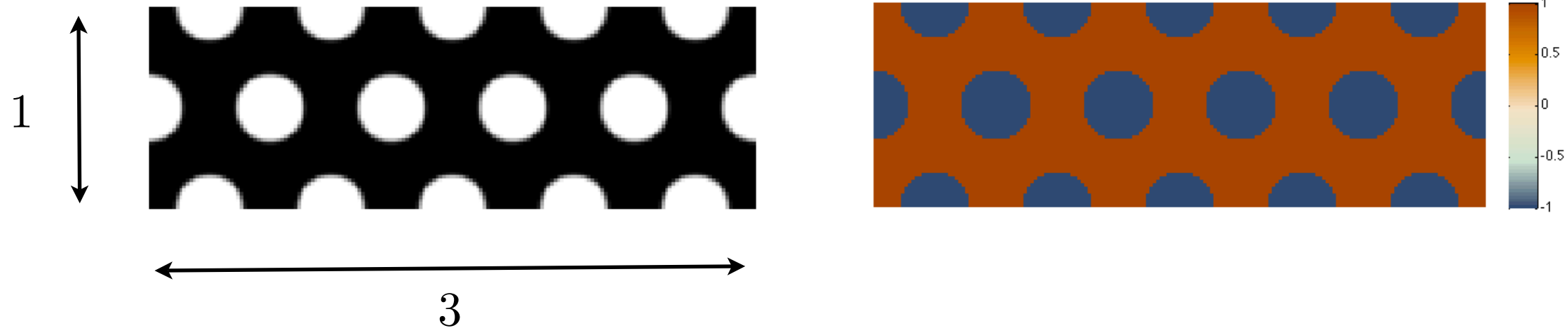
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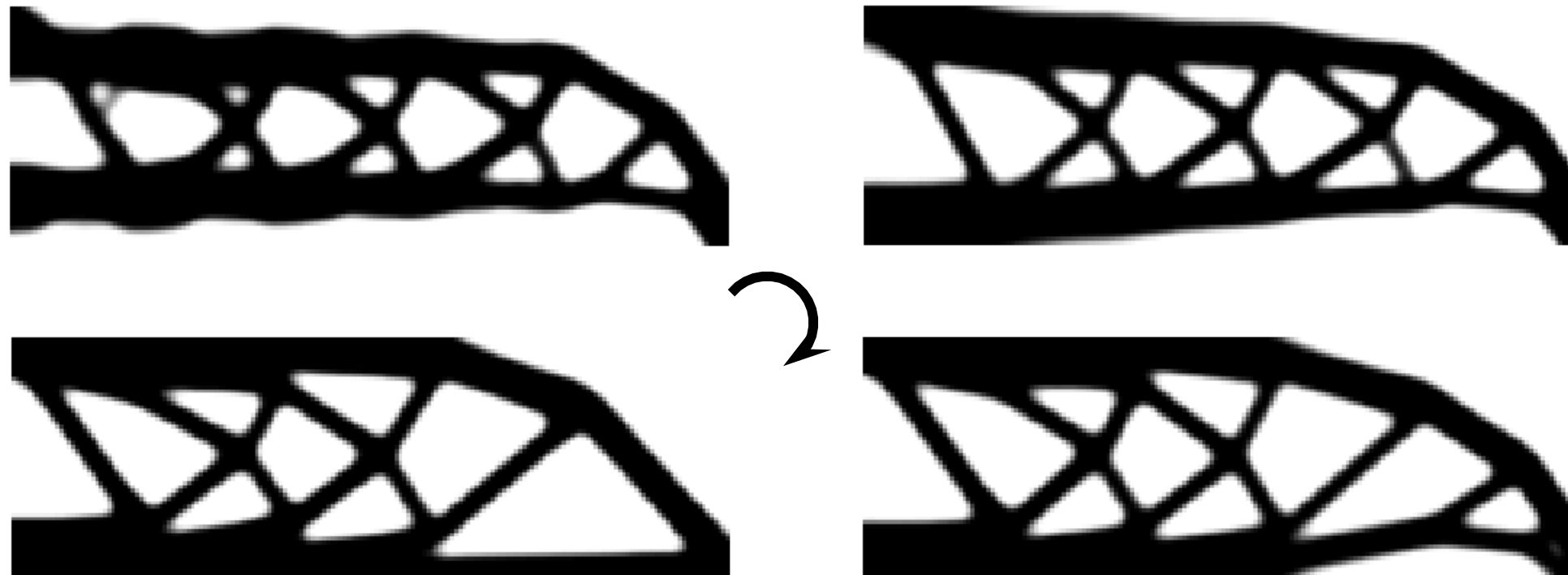
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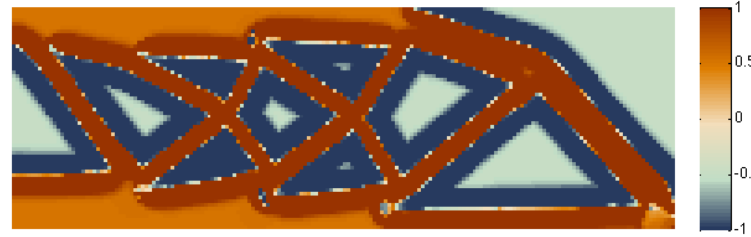
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$$H_w(\varphi)$$



$$\eta$$



$$R = 0.075$$

$$\varphi = K_R \star \eta$$



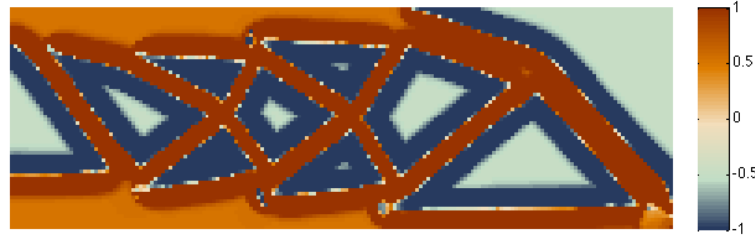
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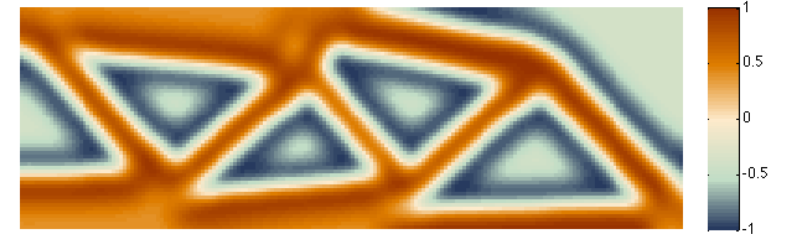
$$\eta$$



$$\varphi = K_R \star \eta$$



$$R = 0.075$$



$$R = 0.100$$

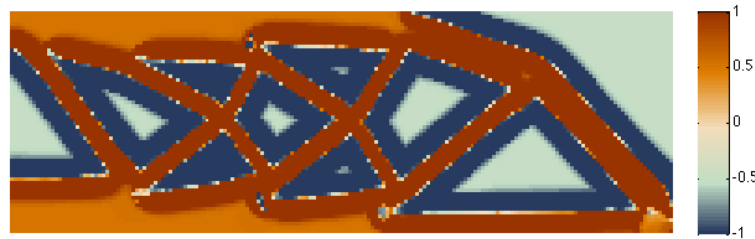
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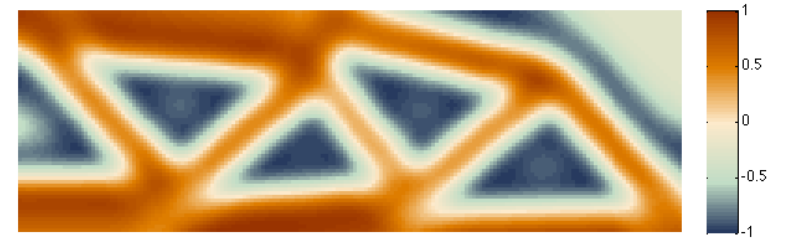
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$$R = 0.150$$

Some numerical results



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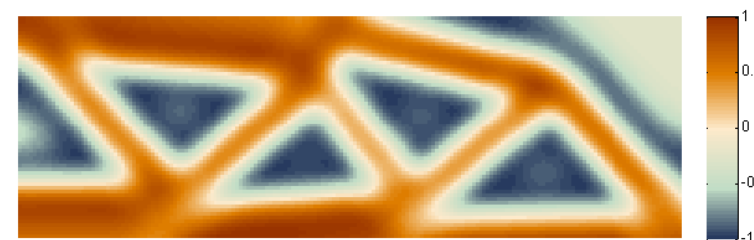
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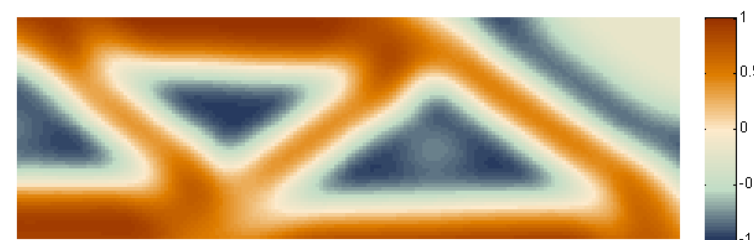
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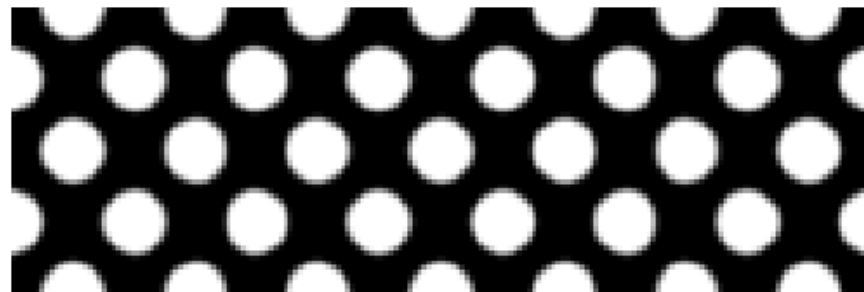


$$R = 0.150$$



$$R = 0.200$$

Initial guess



Final solution





- The nature of the continuum optimal shape problem has implications for the numerical formulations and algorithms
- In addition to smoothness, a uniform “transversality” condition must be imposed on the implicit function field
- Within the restriction framework, the Ersatz material model (filling the voids with compliant material C^-) can be justified

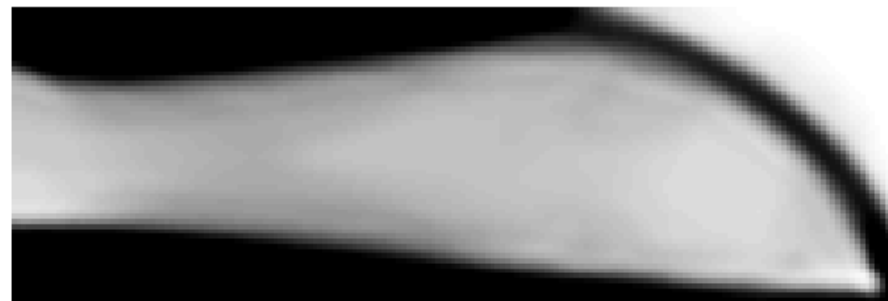
QUESTIONS?



- This fact can be illustrated numerically:



Initial φ



Final solution

- With transversality condition (R2) imposed, however, we can prove that as $w \rightarrow 0$, the optimal solution $\chi_w^* = H_w(\varphi^*)$ converge to solution of problem with $\mathcal{A} = H(\mathcal{F})$