



# Assessment of cohesive traction-separation relationships in ABAQUS: A comparative study



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## ABSTRACT

The definition of a traction-separation relationship is essential in cohesive zone models because it describes the nonlinear fracture process zone. A few models are investigated in this paper and a comparative study is conducted. Among various traction-separation relationships, the one in Abaqus is assessed by evaluating the cohesive traction and its tangent stiffness according to a given separation path. The results demonstrate that the traction-separation relationship in Abaqus can lead to non-physical responses because of a pathological positive tangent stiffness under softening condition. This is reflected in cohesive tractions that increase and decrease repeatedly while the cohesive separation monotonically increases. Thus, together with supporting information, this paper conveys the message that a traction-separation relationship should be developed and selected with great caution, especially under mixed-mode conditions.

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## 1. Introduction

Cohesive zone models have been extensively utilized to represent various nonlinear material fracture behaviors. In many instances, cracks propagate under mixed-mode conditions rather than, for example, isolated mode-I or mode-II. In fact, different choices of the traction-separation relationships can result in different failure behavior and structural response [4,6]. Thus, consistent representation of traction-separation relationships under mixed-mode conditions is critical for proper use and interpretation of cohesive zone models.

In general, traction-separation relationships can be classified into *potential-based models* and *non-potential-based models*. Potential-based models use the concept of cohesive energy potential – see, for example, Needleman [15]. The gradient of the cohesive energy potential leads to the cohesive traction versus separation relationships within the softening region. Several potential expressions have been proposed in conjunction with fracture parameters such as fracture energy and cohesive strength, e.g. [16,28,30]. Following the framework of potential-based models, the so-called PPR (Park-Paulino-Roesler) model was established on the basis of consistent cohesive fracture boundary conditions under mixed-mode conditions [18]. Since its publication, this model has been investi-

gated in several physical contexts including multiscale fracture [5], fragmentation [22,11,24], and composite material failure [17]. The PPR model has also been independently implemented in Abaqus as user-defined elements [20,25]. For an in-depth review of potential-based models, the reader is referred to Park and Paulino [21].

For non-potential-based models, traction-separation models with various shapes have been developed, e.g. linear softening [3], trapezoidal shape [31], bilinear softening [29], trilinear softening [19], and exponential models [12] to name a few. These models do not need to satisfy the symmetry condition, i.e., an exact differential with  $\partial T_n / \Delta_t = \partial T_t / \Delta_n$ . However, non-potential-based models may not provide a consistent traction-separation relationship, and they do not account for all possible separation paths within the softening region.

The traction-separation model proposed by Camanho et al. [3] has been extensively utilized to investigate mixed-mode fracture behavior, e.g. [10,14,23,32]. This model was implemented in the commercial software Abaqus [1] and was reformulated on the basis of cohesive energy potential and damage mechanics [26]. A similar traction-separation relationship was also developed for analyzing delamination of composite materials [9,8]. Recently, McGarry et al. [12] proposed three non-potential-based models while addressing the relative weight of mode-mixity [13]. However, such traction-separation relationships are prone to inconsistency, for instance, a positive stiffness under softening conditions.

In this context, the consistency of the non-potential-based models, including the aforementioned Abaqus model, is assessed, while

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### Nomenclature

$T_n, T_t$	Normal and tangential tractions
$K_p$	Penalty stiffness
$d$	Scalar damage variable
$m, n$	Characteristic parameters in the PPR model
$\Delta_n, \Delta_t$	Normal and tangential separations
$\bar{\Delta}$	Effective separation
$\Gamma_n, \Gamma_t$	Characteristic parameters in the PPR model
$\Psi$	Potential for cohesive fracture
$\alpha, \beta$	Shape parameters in the PPR model
$\alpha_m, \beta_m$	Parameters for weight of the mode-mixity
$\alpha_r$	Parameter for the rate of damage evolution
$\beta_e$	Ratio of tangential separation to normal separation
$\eta$	Parameter for variation of fracture toughness
$\bar{\delta}_c$	Effective critical separation
$\bar{\delta}_f$	Effective complete failure separation
$\delta_n, \delta_t$	Characteristic length scale parameters
$\delta_{nc}, \delta_{tc}$	Characteristic lengths for critical separations
$\lambda_n, \lambda_t$	Initial slope indicators
$\sigma_{max}, \tau_{max}$	Normal and tangential cohesive strengths
$\phi_n, \phi_t$	Normal and tangential fracture energies

several limitations of the models are addressed through comparison with the PPR reference model [18]. The remainder of the paper is organized as follows. Section 2 briefly describes five traction-separation relations, i.e. the Abaqus model, the three models by McGarry et al. [12], and the PPR model. In Section 3, the consistency of the models is verified by evaluating the tangent stiffness of the traction-separation relation. Finally, some key findings of the paper are summarized in Section 4.

## 2. Cohesive zone model

In the cohesive zone model, the nonlinear fracture process zone is approximated by means of cohesive surface elements that follow a traction-separation relation, which are inserted along fracture surfaces. In the commercial software Abaqus, the cohesive elements are inserted a priori between bulk elements (often along a predefined path), leading to the so-called intrinsic cohesive zone model. The intrinsic traction-separation relationships of Abaqus, McGarry et al. and PPR models are briefly explained below.

### 2.1. Traction-Separation relationship in ABAQUS

The traction-separation relationship in Abaqus was originally developed by Camanho et al. [3]. This model considers the variation of fracture toughness as a function of a mode-mixity ratio. In order to account for the mixed-mode condition, an effective separation ( $\bar{\Delta}$ ) is defined as

$$\bar{\Delta} = \sqrt{\Delta_n^2 + \Delta_t^2} \quad (1)$$

where  $\Delta_n$  and  $\Delta_t$  are the normal and tangential separations, respectively. When the effective separation is smaller than an effective critical separation ( $\bar{\delta}_c$ ), the normal and tangential tractions ( $T_n, T_t$ ) are proportional to the normal and tangential separations, which are given as

$$T_n = K_p \Delta_n, \quad T_t = K_p \Delta_t \quad (2)$$

where  $K_p$  is a penalty stiffness. When the separations reach the effective critical separation, the state of separation corresponds to the onset of damage and/or crack initiation. A quadratic failure cri-

terion is employed to determine the onset of damage, which leads to the following effective critical separation, i.e.

$$\bar{\delta}_c = \delta_{nc} \delta_{tc} \sqrt{\frac{1 + \beta_e^2}{\delta_{tc}^2 + (\beta_e \delta_{nc})^2}} \quad (3)$$

where  $\delta_{nc}$  and  $\delta_{tc}$  are defined as the ratio between the normal and tangential cohesive strengths ( $\sigma_{max}, \tau_{max}$ ) and the penalty stiffness, respectively, and  $\beta_e$  is the ratio between the tangential and normal separations. Note that when the normal cohesive strength is the same as the tangential cohesive strength, the critical separation is a constant, i.e.  $\bar{\delta}_c = \sigma_{max}/K_p$ , because the  $\beta_e$  term cancels out.

When the effective separation is greater than the effective critical separation ( $\bar{\delta}_c$ ) and smaller than the effective complete failure separation ( $\bar{\delta}_f$ ), i.e.  $\bar{\delta}_c < \bar{\Delta} < \bar{\delta}_f$ , the state of separation corresponds to the softening condition. Then, the normal and tangential cohesive tractions are defined as

$$T_n = (1 - d) K_p \Delta_n$$

$$T_t = (1 - d) K_p \Delta_t \quad (4)$$

where  $d$  is a scalar damage variable. Two types of damage evolution are available in Abaqus: linear and exponential models. Accordingly, the expressions for the linear and exponential damage models are given as

$$d = \frac{\bar{\delta}_f (\bar{\Delta} - \bar{\delta}_c)}{\bar{\Delta} (\bar{\delta}_f - \bar{\delta}_c)} \quad (5)$$

and

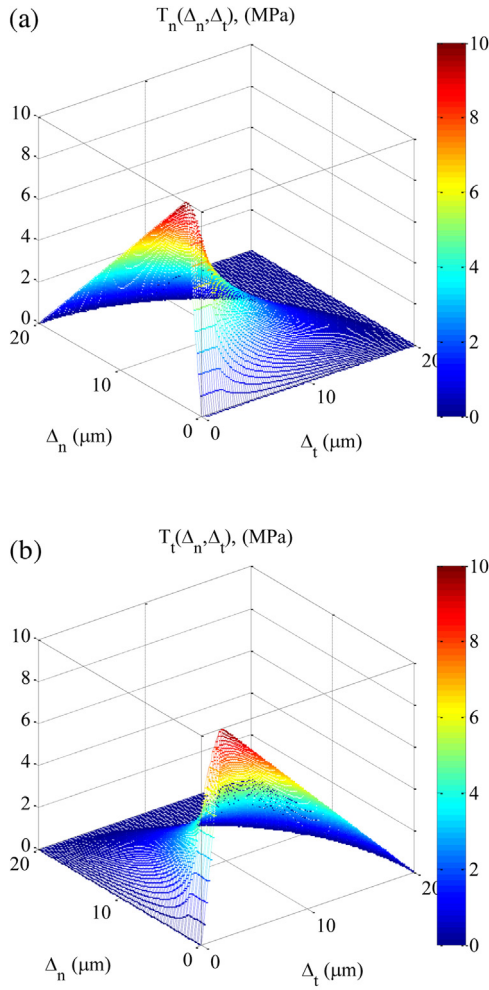
$$d = 1 - \left( \frac{\bar{\delta}_c}{\bar{\Delta}} \right) \left[ 1 - \frac{1 - \exp\left(-\alpha_r \left( \frac{\bar{\Delta} - \bar{\delta}_c}{\bar{\delta}_f - \bar{\delta}_c} \right)\right)}{1 - \exp(-\alpha_r)} \right] \quad (6)$$

respectively, where  $\alpha_r$  is a parameter associated with the rate of damage evolution. The damage variable is initially zero, and monotonically increases to the numeric value of one under softening condition. Additionally, the effective complete failure separation is evaluated by employing the mixed-mode failure criterion proposed by Benzeggagh and Kenane [2], which leads to the following expression:

$$\bar{\delta}_f = \frac{2}{K_p \bar{\delta}_c} \left[ \phi_n + (\phi_t - \phi_n) \left( \frac{\beta_e^2}{1 + \beta_e^2} \right)^\eta \right] \quad (7)$$

where  $\eta$  is a material parameter associated with the variation of fracture toughness. When the normal and tangential fracture energies are the same, the effect of  $\eta$  disappears, and the complete failure separation reduces to  $\bar{\delta}_f = 2\phi_n / (K_p \bar{\delta}_c)$ . If the effective separation is greater than the effective complete failure separation, then both normal and tangential cohesive tractions are zero, as expected. Notice that when the normal and tangential fracture energies and cohesive strengths are the same, the traction-separation relationship in Abaqus simplifies to the model proposed by Tvergaard [27].

To illustrate the aforementioned description, the normal and tangential tractions of the Abaqus model are plotted in Fig. 1. The linear damage model is used, and thus the normal and tangential cohesive tractions linearly decrease along the mode-I and mode-II directions. For the parametric study to be conducted, Mode-I fracture parameters are arbitrarily selected as the fracture energy of 100 J/m<sup>2</sup> and the cohesive strength of 10 MPa. The mode-II fracture parameters are assumed to be the same as the mode-I parameters, leading to constant values of the effective critical separation and the complete failure separation with respect to the mode-mixity. In addition, the penalty stiffness is selected as 5 MPa/ $\mu$ m, which



**Fig. 1.** Illustration of the Abaqus CZM: (a) Normal and (b) tangential cohesive tractions with  $\phi_n = 100 \text{ J/m}^2$ ,  $\phi_t = 100 \text{ J/m}^2$ ,  $\sigma_{max} = 10 \text{ MPa}$ ,  $\tau_{max} = 10 \text{ MPa}$ , and  $K_p = 5 \text{ MPa}/\mu\text{m}$ .

results in  $\bar{\delta}_c/\bar{\delta}_f = 0.1$  where  $\bar{\delta}_c$  and  $\bar{\delta}_f$  are  $2 \mu\text{m}$  and  $20 \mu\text{m}$ , respectively, in this example.

## 2.2. Non-potential-based models by McGarry et al. [12]

Based on the exponential potential by Xu and Needleman [30], McGarry et al. [12] proposed three non-potential-based models. The non-potential-based formulation 1 (NP1) is expressed as

$$T_n(\Delta_n, \Delta_t) = \sigma_{max} \exp(1) \left( \frac{\Delta_n}{\delta_n} \right) \exp\left(-\frac{\Delta_n}{\delta_n}\right) \exp\left(-\frac{\Delta_t^2}{\delta_t^2}\right) \quad (8)$$

and

$$T_t(\Delta_n, \Delta_t) = \tau_{max} \sqrt{2 \exp(1)} \left( \frac{\Delta_t}{\delta_t} \right) \exp\left(-\frac{\Delta_n}{\delta_n}\right) \exp\left(-\frac{\Delta_t^2}{\delta_t^2}\right) \quad (9)$$

where the characteristic length ( $\delta_n, \delta_t$ ) is evaluated from the fracture energy. In order to provide equivalent normal and tangential traction-separation relationships, the second non-potential-based formulation (NP2) is given as

$$T_n(\Delta_n, \Delta_t) = \sigma_{max} \exp(1) \left( \frac{\Delta_n}{\delta_n} \right) \exp\left(-\frac{\Delta_n}{\delta_n}\right) \exp\left(-\alpha_m \sqrt{\frac{\Delta_t^2}{\delta_t^2}}\right) \quad (10)$$

and

$$T_t(\Delta_n, \Delta_t) = \tau_{max} \exp(1) \left( \frac{\Delta_t}{\delta_t} \right) \exp\left(-\sqrt{\frac{\Delta_n^2}{\delta_n^2}}\right) \exp\left(-\beta_m \frac{\Delta_n}{\delta_n}\right) \quad (11)$$

where  $\alpha_m$  and  $\beta_m$  are associated with the weight of the mode-mixity. Finally, based on the effective separation, the third non-potential-based formulation (NP3), also named as separation magnitude coupling (SMC) formulation, is given as

$$T_n(\Delta_n, \Delta_t) = \sigma_{max} \exp(1) \left( \frac{\Delta_n}{\delta_n} \right) \exp\left(-\sqrt{\frac{\Delta_n^2}{\delta_n^2} + \frac{\Delta_t^2}{\delta_t^2}}\right) \quad (12)$$

and

$$T_t(\Delta_n, \Delta_t) = \tau_{max} \exp(1) \left( \frac{\Delta_t}{\delta_t} \right) \exp\left(-\sqrt{\frac{\Delta_n^2}{\delta_n^2} + \frac{\Delta_t^2}{\delta_t^2}}\right) \quad (13)$$

Note that for the case of  $\delta_n = \delta_t$ , the NP3 formulation is a special case of the model proposed by Tvergaard [27]. For further details of the models NP1, NP2 and NP3, the reader is referred to McGarry et al. [12].

## 2.3. PPR model

In order to provide consistent cohesive traction-separation relationships, the PPR model by Park et al. [18] has been formulated on the basis of a fracture energy potential. The potential of the PPR model is given as

$$\Psi(\Delta_n, \Delta_t) = \min(\phi_n, \phi_t) + \left[ \Gamma_n \left(1 - \frac{\Delta_n}{\delta_n}\right)^\alpha \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^m + \langle \phi_n - \phi_t \rangle \right] \left[ \Gamma_t \left(1 - \frac{|\Delta_t|}{\delta_t}\right)^\beta \left(\frac{n}{\beta} + \frac{|\Delta_t|}{\delta_t}\right)^n + \langle \phi_t - \phi_n \rangle \right] \quad (14)$$

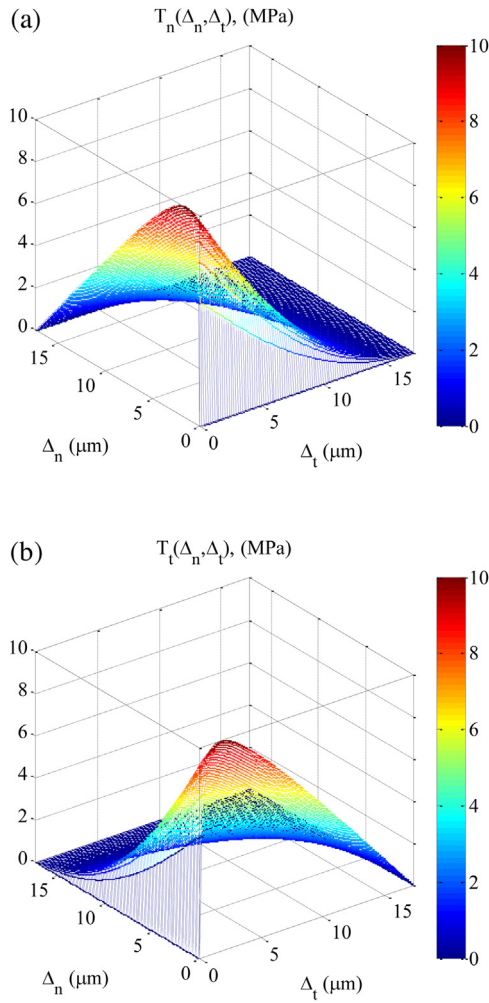
where  $\langle \cdot \rangle$  is the Macaulay bracket. Characteristic parameters ( $\Gamma_n, \Gamma_t; \delta_n, \delta_t; m, n; \alpha, \beta$ ) are explicitly obtained from actual cohesive fracture parameters such as fracture energy, cohesive strength, shape of interactions, and initial ascending slope for each fracture mode. The expressions for the characteristic parameters are given in the original paper by Park et al. [18]. The gradient of the potential leads to the normal and tangential tractions, expressed as

$$T_n(\Delta_n, \Delta_t) = \frac{\Gamma_n}{\delta_n} \left[ m \left(1 - \frac{\Delta_n}{\delta_n}\right)^\alpha \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^{m-1} - \alpha \left(1 - \frac{\Delta_n}{\delta_n}\right)^{\alpha-1} \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^m \right] \times \left[ \Gamma_t \left(1 - \frac{|\Delta_t|}{\delta_t}\right)^\beta \left(\frac{n}{\beta} + \frac{|\Delta_t|}{\delta_t}\right)^n + \langle \phi_t - \phi_n \rangle \right] \quad (15)$$

and

$$T_t(\Delta_n, \Delta_t) = \frac{\Gamma_t}{\delta_t} \left[ n \left(1 - \frac{|\Delta_t|}{\delta_t}\right)^\beta \left(\frac{n}{\beta} + \frac{|\Delta_t|}{\delta_t}\right)^{n-1} - \beta \left(1 - \frac{|\Delta_t|}{\delta_t}\right)^{\beta-1} \left(\frac{n}{\beta} + \frac{|\Delta_t|}{\delta_t}\right)^n \right] \times \left[ \Gamma_n \left(1 - \frac{\Delta_n}{\delta_n}\right)^\alpha \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^m + \langle \phi_t - \phi_n \rangle \right]. \quad (16)$$

In general, the cohesive tractions increase from zero to their maximum value when the cohesive separations go from zero to the critical separations ( $\delta_{nc}, \delta_{tc}$ ). Then, the tractions monotonically decrease to zero in a softening region when the cohesive separations reach complete failure separations ( $\delta_n, \delta_t$ ). Note that the cohesive tractions (Eqs. (15) and (16)) are defined within the cohesive interaction region, which is associated with the complete failure conditions. If the separation is outside of the cohesive interaction region, the cohesive tractions are zero. Additionally, the



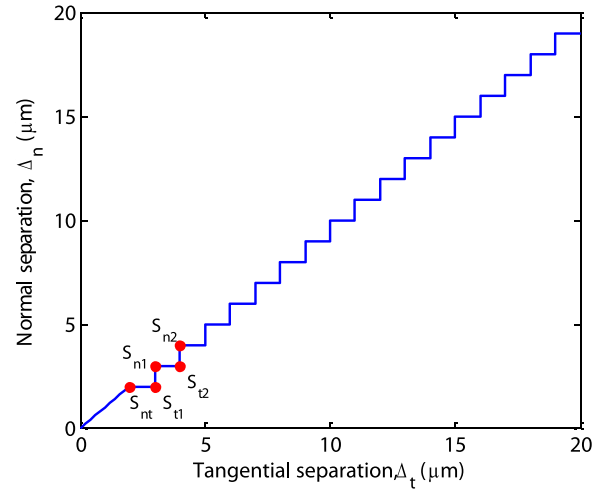
**Fig. 2.** PPR Model: (a) Normal and (b) tangential cohesive tractions with  $\phi_n = 100 \text{ J/m}^2$ ,  $\phi_t = 100 \text{ J/m}^2$ ,  $\sigma_{max} = 10 \text{ MPa}$ ,  $\tau_{max} = 10 \text{ MPa}$ ,  $\alpha = 2$ ,  $\beta = 2$ ,  $\lambda_n = 0.1$ , and  $\lambda_t = 0.1$ .

cohesive interaction region consists of two parts, i.e. elastic region and softening region. The boundary between the elastic region and the softening region is defined by the critical separation.

To illustrate, the normal and tangential tractions of the PPR model are plotted in Fig. 2 within the cohesive interaction region. For mode-I, the fracture energy ( $\phi_n$ ) and the cohesive strength ( $\sigma_{max}$ ) are  $100 \text{ J/m}^2$  and  $10 \text{ MPa}$ , respectively. The initial slope indicator ( $\lambda_n$ ), i.e. the ratio of the critical separation ( $\delta_{nc}$ ) to the complete separation ( $\delta_n$ ), is selected as 0.1. The shape parameter ( $\alpha$ ) is 2, which provides nearly linear softening behavior. The fracture parameters of the tangential mode ( $\phi_t$ ,  $\tau_{max}$ ,  $\lambda_t$ ,  $\beta$ ) are assumed to be the same as the fracture parameters of mode-I.

### 3. On consistency of traction-separation relationships

One of the fundamental requirements for a consistent traction-separation relationship is that the cohesive traction decreases according to the increase of the separation across fracture surfaces, which results in a negative tangent stiffness under the softening condition [7,21]. This is because the wider the separation is, the more the damage that occurs, and thus less resistance is expected along the fracture surfaces. In order to assess the consistency of the traction-separation relationships of Abaqus, NP1, NP2, NP3,



**Fig. 3.** Arbitrary separation path for evaluating normal and tangential cohesive tractions of various cohesive zone models.

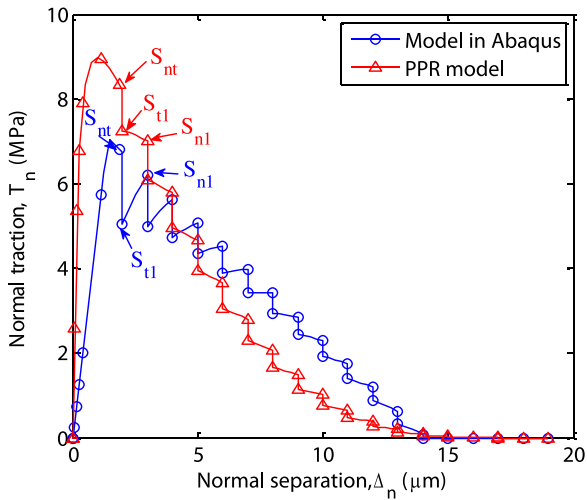
and PPR models, the cohesive traction is evaluated for an arbitrary separation path. Then, whenever non-physical behavior occurs, it is explained by calculating the tangent stiffness of the traction-separation relations under the softening condition. Finally, the onset of softening is evaluated to demonstrate the consistency of the traction-separation relationships.

#### 3.1. Traction-separation relationship for an arbitrary separation path

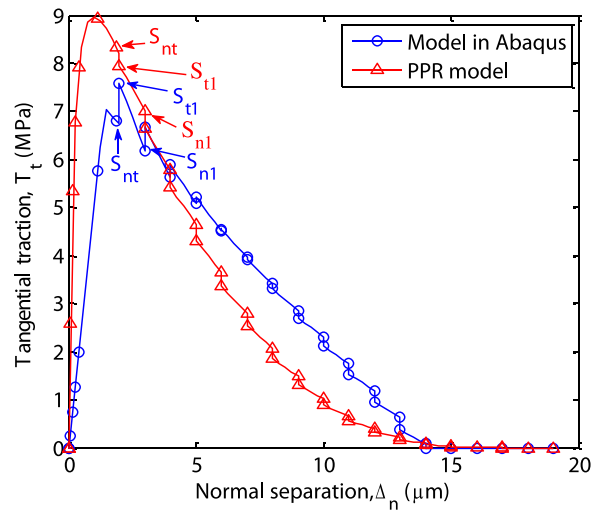
An arbitrary separation path, as shown in Fig. 3, is selected to evaluate the aforementioned cohesive zone models (discussed in Section 2). The cohesive fracture parameters used to assess the cohesive zone models are as described in Figs. 1 and 2. For the traction-separation relationships, the fracture energy and the cohesive strength are  $100 \text{ J/m}^2$  and  $10 \text{ MPa}$ , respectively, and the ratio of the critical separation to the complete failure separation is 0.1. Additionally, the critical separations of the Abaqus and PPR models are  $2 \mu\text{m}$  and  $1.732 \mu\text{m}$ , respectively, while of the NP1, NP2 and NP3 models are not clearly defined.

For the separation path shown in Fig. 3, both normal and tangential separations initially increase with a magnitude of  $2 \mu\text{m}$ , which corresponds to Point  $S_{nt}$ . Next, the tangential separation increases an amount of  $1 \mu\text{m}$  (Point  $S_{t1}$ ), and then the normal separation increases by  $1 \mu\text{m}$  (Point  $S_{n1}$ ). Afterwards, this step is repeated until the state of separation reaches the complete failure condition. Because the specified separation monotonically increases, one expects a monotonic increase of damage, resulting on a monotonic decrease of the cohesive traction within the softening region. If a separation path is within the elastic region, i.e. the initial ascending part, the cohesive traction can increase while the separation increases, as expected.

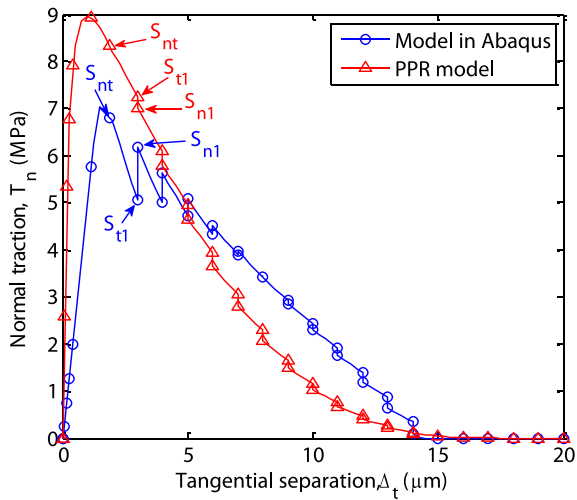
Based on the specified separation path, the normal and tangential cohesive tractions are plotted in Figs. 4 and 5 for both the Abaqus and the PPR models. The normal and tangential cohesive tractions initially increase up to the critical separation, and then monotonically decrease until normal and tangential separations reach Point  $S_{nt}$  for both models. This is because the state of separation at Point  $S_{nt}$  corresponds to the softening condition. Note that, for the Abaqus model, the effective critical separation ( $\bar{\Delta}$ ) is  $2.828 \mu\text{m}$ , which is greater than the effective critical separation (i.e.  $\bar{\delta}_c = 2 \mu\text{m}$ ). For the PPR model, the normal and tangential separations (i.e.  $\Delta_n = \Delta_t = 2 \mu\text{m}$ ) are also greater than the critical normal and tangential separations (i.e.  $\delta_{nc} = \delta_{tc} = 1.732 \mu\text{m}$ ).



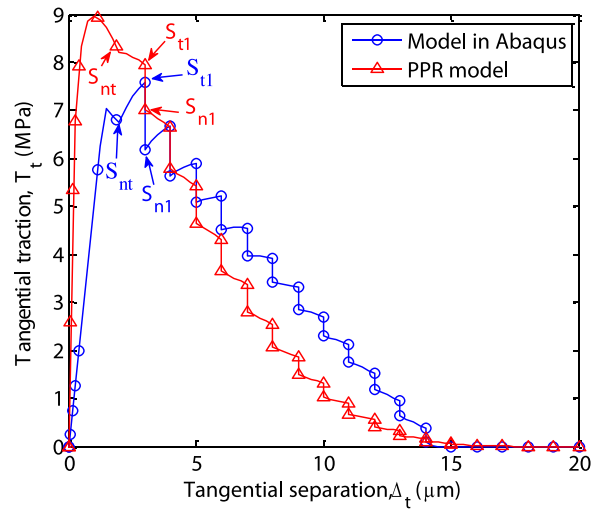
(a)



(a)



(b)



(b)

**Fig. 4.** Normal cohesive traction with respect to (a) normal separation and (b) tangential separation.

After reaching Point  $S_{nt}$ , the cohesive traction response of the Abaqus model is different from the response of the PPR model. For the Abaqus model, when the tangential separation increases from  $2 \mu\text{m}$  to  $3 \mu\text{m}$  (from Point  $S_{nt}$  to Point  $S_{t1}$ ), the normal traction decreases (see Fig. 4(b)) while the tangential traction increases (see Fig. 5(b)). Note that the increase of the cohesive traction with respect to the increase of separation is not realistic because a larger separation should be associated with more damage and less resistance along the fracture surfaces. Next, when the normal separation increases from  $2 \mu\text{m}$  to  $3 \mu\text{m}$  (from Point  $S_{t1}$  to Point  $S_{n1}$ ), the normal traction increases (see Fig. 4(a)) and the tangential traction decreases (see Fig. 5(a)), in the Abaqus model. These cohesive traction fluctuations are non-physical, and persist until the normal and tangential separations approximately reach  $6.75 \mu\text{m}$  and  $7.25 \mu\text{m}$ , respectively. When the normal and tangential separations are greater than these values, then both normal and tangential tractions decrease (in this example). However, for the PPR model, both normal and tangential separations monotonically decrease within the entire softening region.

The oscillation of the cohesive tractions in the Abaqus model is associated with the competition between the damage variable and a linear separation term. The first term in Eq. (4), i.e.

**Fig. 5.** Tangential cohesive traction with respect to (a) normal separation and (b) tangential separation.

(1–d), contributes to the decrease of the cohesive traction because the damage variable ( $d$ ) monotonically increases according to the increase of the separation variable. The linear separation terms in Eq (4), i.e.  $\Delta_n$  and  $\Delta_t$ , contribute to the increase of the cohesive traction. Then, when the rate of change of the linear separation term is greater than the rate of the change of the damage variable term, the cohesive traction increases with the increase of the cohesive separation. In order to confirm such behavior, the tangent stiffness under the softening region is evaluated in the following subsection.

Additionally, the cohesive traction-separations of the non-potential-based models, i.e. NP1, NP2 and NP3, are evaluated for the same separation path, as shown in Fig. 6. Similarly to the Abaqus model, the non-potential-based models demonstrate the oscillation of the cohesive traction according to the increase of separation. Note that the oscillation is also associated with the competition between the linear separation term (i.e.  $\Delta_n$  and  $\Delta_t$ ) and the exponential decaying terms in the non-potential-based models.



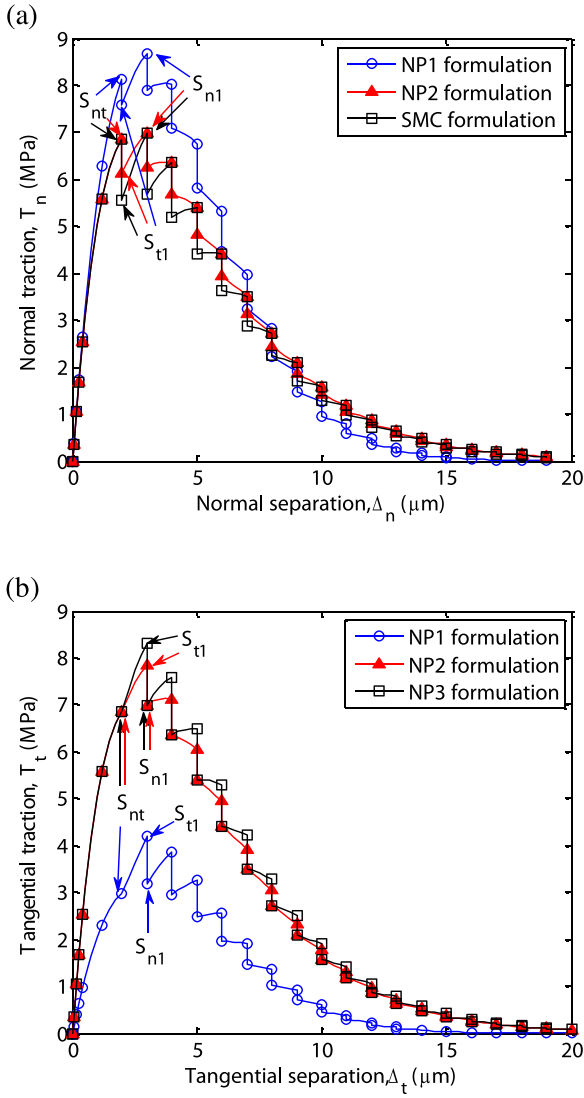


Fig. 6. Cohesive traction of the non-potential-based models: (a) normal traction and (b) tangential traction.

### 3.2. Tangent stiffness under the softening region

In order to confirm the non-physical behavior of the Abaqus model, the tangent stiffness is evaluated in the softening region. For this purpose, we employ the chain rule to obtain the derivative of the damage variable from Eq. (5) with respect to the normal separation, i.e.

$$\frac{\partial d}{\partial \Delta_n} = \frac{\partial d}{\partial \bar{\Delta}} \frac{\partial \bar{\Delta}}{\partial \Delta_n} + \frac{\partial d}{\partial \bar{\delta}_c} \frac{\partial \bar{\delta}_c}{\partial \Delta_n} + \frac{\partial d}{\partial \bar{\delta}_f} \frac{\partial \bar{\delta}_f}{\partial \Delta_n} \quad (17)$$

because  $\bar{\Delta}$ ,  $\bar{\delta}_c$  and  $\bar{\delta}_f$  are expressed in terms of the normal and tangential separations. Alternatively, Turon et al. [26] assumed that the variations of the critical separation and the complete failure separation were not significant, and thus only considered the first term in Eq. (17) for the derivative of the damage variable. Using Eq. (17), the derivative of the normal traction (Eq. (4)) with respect to the normal separation is given as

$$\frac{\partial T_n}{\partial \Delta_n} = \frac{\Gamma_n}{\delta_n^2} \left[ (m^2 - m) \left(1 - \frac{\Delta_n}{\delta_n}\right)^\alpha \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^{m-2} + (\alpha - \alpha^2) \left(1 - \frac{\Delta_n}{\delta_n}\right)^{\alpha-2} \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^m - 2\alpha m \left(1 - \frac{\Delta_n}{\delta_n}\right)^{\alpha-1} \left(\frac{m}{\alpha} + \frac{\Delta_n}{\delta_n}\right)^m \right] \left[ \Gamma_t \left(1 - \frac{|\Delta_t|}{\delta_t}\right)^\beta \left(\frac{n}{\beta} + \frac{|\Delta_t|}{\delta_t}\right)^n + (\phi_t - \phi_n) \right] \quad (19)$$

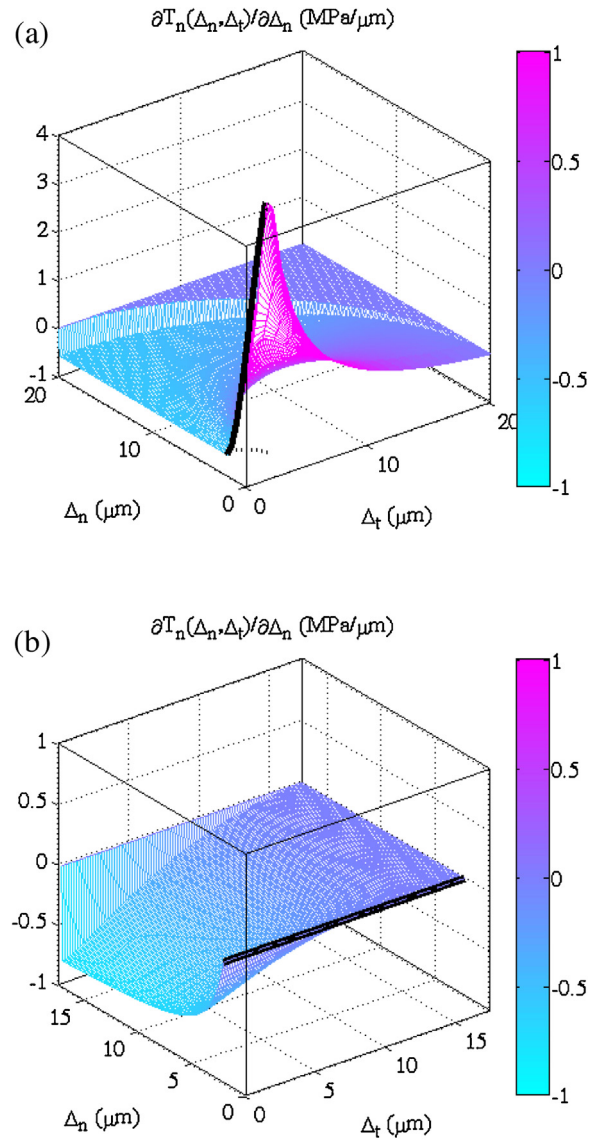


Fig. 7. Normal material tangent stiffness of (a) Abaqus model and (b) PPR model.

$$\frac{\partial T_n}{\partial \Delta_n} = K_p (1 - d) - \frac{K_p \Delta_n^2 \bar{\delta}_c \bar{\delta}_f}{\bar{\Delta}^3 (\bar{\delta}_f - \bar{\delta}_c)} \left( 1 + \frac{\Delta_t^2 (\bar{\delta}_c + \bar{\delta}_f - 2\bar{\Delta})}{\bar{\Delta}^2 (\bar{\delta}_f - \bar{\delta}_c)} \left( \frac{\bar{\delta}_c^2}{\delta_{nc}^2} - \frac{\bar{\delta}_c^2}{\delta_{tc}^2} \right) \right) - \frac{2\eta d \Delta_n^2 \bar{\delta}_c}{\bar{\Delta}^2 (\bar{\delta}_f - \bar{\delta}_c)} \left( K_p - \frac{2\phi_n}{\bar{\delta}_c \bar{\delta}_f} \right) \quad (18)$$

within the softening condition, i.e.  $\bar{\delta}_c < \bar{\Delta} < \bar{\delta}_f$ .

For the PPR model, the complete failure separation and the characteristic parameters are obtained from the fracture parameters, and thus the derivative of the normal traction with respect to the normal separation is expressed as

Note that for the normal cohesive traction, the onset of the softening condition occurs when the normal separation is greater than the normal critical separation, which corresponds to the softening region.

The normal tangent stiffness within the softening region for each model is plotted in Fig. 7 with the fracture parameters used in Section 3.1. For the Abaqus model (Fig. 7(a)), positive stiffness is observed in a portion of the softening region, which results in the increase of the cohesive traction with respect to the increase of the cohesive separation. Because of such positive stiffness, the oscillation of the cohesive tractions shown in Figs. 4 and 5 is expected. However, for the PPR model, Fig. 7(b) clearly demonstrates that the stiffness within the softening region is always negative. This explains why the PPR cohesive tractions monotonically decrease within the softening region, as observed in Figs. 4 and 5.

The tangent stiffness along the critical separation is indicated by a thick solid line (see Fig. 7). For the model in Abaqus, before reaching the onset of crack initiation, the tangent stiffness is initially constant, i.e. the penalty stiffness (see Eq. (2)). When the effective separation reaches the critical separation, the magnitude and sign of the tangent stiffness change according to the ratio of the normal separation to the tangential separation. Thus, a discontinuous tangent stiffness within the cohesive interaction region is expected. However, the PPR model provides zero stiffness along the critical separation, which results in the continuous stiffness distribution within the cohesive interaction region.

### 3.3. Onset of softening for an arbitrary separation path

In order to check the consistency of the Abaqus, PPR, and non-potential based models, the onset of softening is evaluated for an arbitrary separation path. The normal separation initially increases up to  $0.5 \mu\text{m}$  (denoted as **Point S<sub>A</sub>**), and then the tangential separation increase up to  $1 \mu\text{m}$  (denoted as **Point S<sub>B</sub>**). Afterwards, the normal separation increases up to the complete failure condition. For the given separation path, the traction-separation relationships of the Abaqus, PPR, and the non-potential based models result in similar cohesive traction variations but different positions of the onset of softening.

As shown in Fig. 8, for the Abaqus, PPR, and non-potential based models, the normal traction initially increases within the elastic region while the normal separation increases. As the tangential separation increases from **Point S<sub>A</sub>** to **Point S<sub>B</sub>**, the normal traction decreases, because the increase of tangential separation results in the decrease of the cohesive strength under mixed-mode condition. Then, while the normal separation increases, the cohesive traction reaches a maximum value and monotonically decreases to zero. The corresponding cohesive tractions of Abaqus, PPR and NP3 models are plotted using a black solid line, while the cohesive traction along the critical separation is indicated as a red dashed line within the normal traction-separation relationships, as shown in Fig. 8. Note that the NP3 model is selected as a representative example of the non-potential based models by McGarry et al. [12]. The cohesive traction along the critical separation is not plotted for the NP3 model (see Fig. 8(c)) because the onset of softening is not clearly defined in this non-potential based model.

In the Abaqus model, the onset of softening is reached when the tangential separation increases from **Point S<sub>A</sub>** to **Point S<sub>B</sub>**, which corresponds to the intersection of the red dashed line and the black solid line in Fig. 8(a). After the onset of softening is reached, the cohesive normal traction predicted by the Abaqus model increases as the normal separation increases, as discussed previously.

In contrast to the Abaqus model, the onset of softening of the PPR model corresponds to the maximum traction, as shown by the intersection between the red dashed line and the black solid line in Fig. 8(b). This leads to a decrease of the cohesive traction

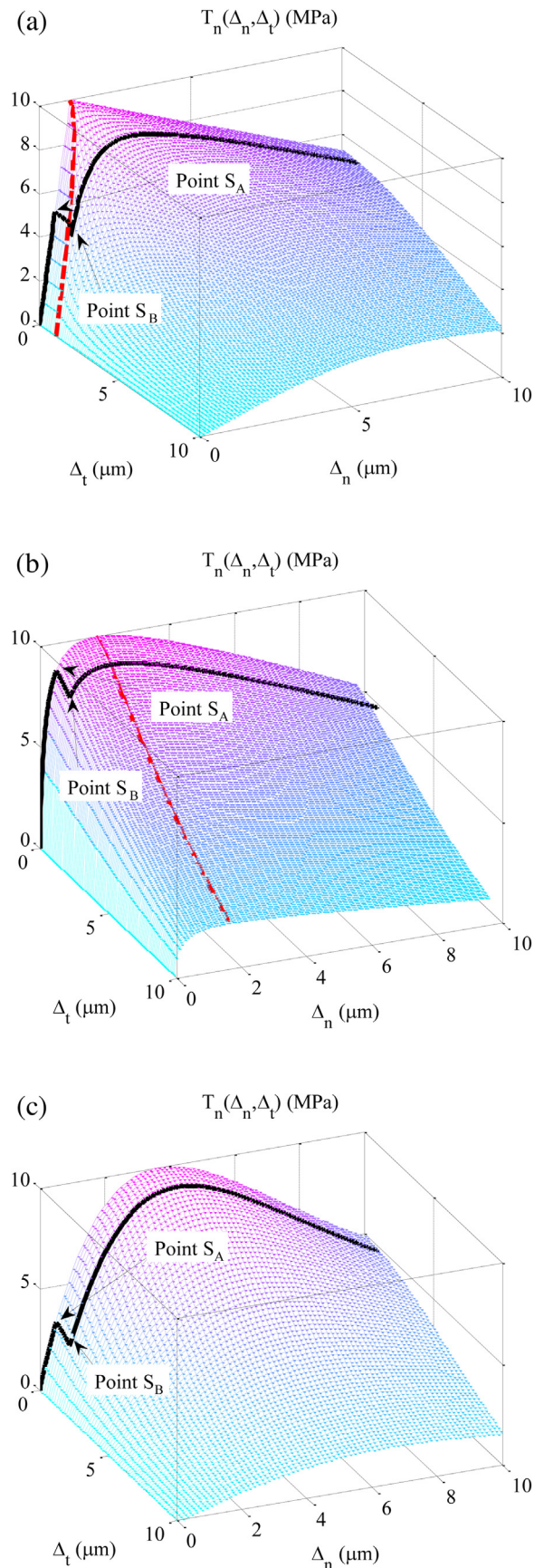


Fig. 8. Cohesive traction for a given separation path within the normal traction-separation relationships (a) Abaqus model, (b) PPR model, and (c) NP3 model.

as the normal separation increases in the softening region, thus demonstrating consistency of the PPR model.

For the non-potential NP3 model, the onset of softening is not clearly defined in the traction-separation relationship, and thus it may be assumed on the basis of: (1) the sign of the rate of cohesive traction, or (2) the maximum cohesive traction. First, the onset of softening can be assumed when the sign of the rate of cohesive traction changes from positive to negative, which corresponds to **Point S<sub>A</sub>** in this example. After reaching **Point S<sub>B</sub>**, the cohesive traction of the NP3 model increases up to 9.65 MPa while the cohesive traction at **Point S<sub>A</sub>** is 3.23 MPa. Therefore, this assumption results in the onset of softening at a very low cohesive traction, which is not physical.

Alternatively, the onset of softening can be defined as a maximum cohesive traction in a separation history. However, in this case, one cannot identify the onset of softening before knowing the entire separation path, which is not possible in general computational simulations.

#### 4. Conclusions

A comparative study of traction-separation relationships in cohesive zone models has been presented. Specifically, the traction-separation relationship in Abaqus is assessed in order to check the consistency of the constitutive model, and is compared with the traction-separation relationship of the PPR model. The model in Abaqus can provide non-physical behavior for certain separation paths because the model does not always provide a negative tangent stiffness within the softening region. Similar issues were observed with the so-called NP1, NP2 and NP3 models by McGarry et al. [12]. Thus, the cohesive traction can increase and decrease repeatedly while the separation monotonically increases under softening conditions. Although the Abaqus model may provide satisfactory results for certain classes of problems, the oscillation and the positive stiffness in the traction-separation relation can adversely impact instances such as numerical stability, multiple crack interactions, and dissipated energy. In comparison, the PPR model does not display such non-physical behavior and guarantees the negative tangent stiffness within the softening region. In conclusion, traction-separation relationships should be carefully developed and selected for mixed-mode failure investigations.

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